# VELOCITY MEASUREMENTS ON FLOW AROUND A CYLINDER 

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#### Abstract

If the cylinder is set in an open channel, the flow upstream will undergo a separation of the turbulent boundary layer and rolls up to form the well-known horseshoe-vortex system, which is swept around the cylinder. The flow fields around a cylinder positioned normal to the flow in an open channel have been investigated experimentally for two types of flow. An Acoustic Doppler Velocity Profiler was used to obtain instantaneously the three directions of the velocity in the flow. Results of the experiments present the vertical velocities distributions of the flow around a cylinder, showing a highly three-dimensional picture, being reasonably organized in the front and less organized in the back of the cylinder, where flow separation takes place. It also shows that a horseshoe-vortex system is established, existing of a measurable vortex with underneath a return flow. The system develops itself in the upstream corner at the nose of the cylinder and stretches around towards the downstream.


Keywords: flow fields around a cylinder, vertical velocities distributions, horseshoe-vortex

## INTRODUCTION

Frown around an obstacle installed in an open channel, creates a three dimensional flow fields. The flow upstream will undergo a separation of the turbulent boundary layer and rolls up to form the well-known horseshoe-vortex system, which is swept around the cylinder. This type of flow occurs in a variety of situations, such as flow around bridge piers, around buildings and structures, and at different types of junctions.

The experiments done with airflow over round and round-nosed objects have been investigated by numerous authors ( see Baker, 1980, Eckerle and Langston, 1987, Pierce and Tree, 1990, Eckerle and Awad, 1991, Agui and Andreopoulos, 1992, Monti, 1994, Monnier and Stanislas, 1996) and with free surface flow ( see Bolcs, 1969, Dargahi, 1987, Magini, 1993, Graf and Yulistiyanto, 1997). These are limited to the flow in the vertical plane of symmetry, in plane $\alpha=0^{\circ}$. Similar experiment measurements were conducted by Abdul Karim Barbhuiya and Subhasish Dey (2003), for flow around a vertical semicircular cylinder attached to the sidewall and at a wing-wall abutment installed at a rectangular channel, and by Andreas Roulund at.al. (2005).

This study investigates flow fields around a cylinder by using experiment data conducted in open channel (Yulistiyanto, B., 1997).

## RESEARCH METHOD

The velocity distributions in different planes upstream and downstream from a cylinder have been
performed, whose hydraulic characteristics are given in Table 1.

The experiments were performed in a 43.0 [m] long and $2.0[\mathrm{~m}]$ wide tilting flume. The working section of the flume was positioned at $x=16.0$ [m] downstream from the entrance, where a cylinder, having a diameter of $\mathrm{D}=22.0[\mathrm{~cm}]$ and a height of H $=0.5[\mathrm{~m}]$, was installed, being positioned normal to the flow. The uniform approach flow described previously ( see Graf and Yulistiyanto, 1997) was fully developed and may be considered to be twodimensional, $\mathrm{B} / \mathrm{h}_{\infty}>5$, turbulent and subcritical.

An Acoustic Doppler Velocity Profiler (ADVP) was used to measure the velocities. This nonintrusive instrument measures instantaneously the vertical profiles of the velocity in the three directions as well as their turbulence parameters. In this work, the ADVP-instrument operates in bistatic and tristatic mode using two and three transducers, respectively, at a time. The ADVP-instrument, using a measuring frequency of 12 Hz , was installed at a fixed location below the channel bed. This obliged us to move the position of the cylinder, mounted on a moveable carriage, around the ADVP-instrument. Detailed information on the experimental installation and experimental results is found elsewhere (see Yulistiyanto, 1997).

Table 1. Flow variables and channel parameters

| Cylinder |  |  |  | Channel:uniform flow $\mathrm{B}=2.0[\mathrm{~m}]$ |  |  |  |  |  |  |  |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| Test | D <br> $[\mathrm{m}]$ | $\mathrm{Re}_{\mathrm{D}}$ <br> $10^{5}$ | Q <br> $\left[\mathrm{m}^{3} / \mathrm{dt}\right]$ | $\mathrm{S}_{\mathrm{o}}$ <br> $10^{-4}$ | n <br> $\left[\mathrm{m}^{-1 / 3} \mathrm{~s}\right]$ | $\mathrm{U}_{\infty}$ <br> $[\mathrm{m} / \mathrm{dt}]$ | $\mathrm{h}_{\infty}$ <br> $[\mathrm{m}]$ | $\mathrm{B} / \mathrm{h}$ | $\mathrm{U}_{*_{\text {cr }}}$ <br> $[\mathrm{m} / \mathrm{dt}]$ | $\mathrm{Re}_{\mathrm{h}}$ <br> $10^{5}$ | Fr |  |  |  |  |  |
| 1 | 0.22 | 1.48 | 0.248 | 6.25 | 0.012 | 0.67 | 0.185 | 10.8 | 0.029 | 1.24 | 0.5 |  |  |  |  |  |
| 2 | 0.22 | 0.95 | 0.149 | 2.80 | 0.012 | 0.43 | 0.173 | 11.6 | 0.021 | 0.74 | 0.33 |  |  |  |  |  |

## RESULTS AND DISCUSSION

## Measured Velocity Distributions

Velocities distributions, $\bar{u}(z), \bar{v}(z), \bar{w}(z)$, were measures at different points, $\mathrm{Q}(\mathrm{x}, \mathrm{y}, \mathrm{z})$ or $\mathrm{Q}(\mathrm{r}, \alpha, \mathrm{z})$, in the range of $12<\mathrm{r}[\mathrm{cm}]<44$ and $0<\alpha\left[^{\circ}\right]<180$ around the cylinder; the origin was positioned in the centre of the cylinder, $\mathrm{Q}(0,0, \mathrm{z})$, whose radius is $\mathrm{r}=11[\mathrm{~cm}]$.

The vertical distributions of the velocities for test 2 are shown in Figure 1; for test 1 they are rather similar (see Graf and Yulistiyanto, 1998). The data are made dimensionless, using the approach velocity, $\mathrm{U}_{\infty}$, and the local flow depth, h . The flow around the cylinder is three-dimensional, being reasonably organized at the upstream and much less organized at the downstream of the cylinder. The following can be observed:

Approaching the cylinder one observes for the longitudinal velocity profiles, $\bar{u}(z) / U_{\infty}$, the following:

- In the plane $0^{\circ}$, the profiles decrease, having more uniform distribution at value of $\bar{u}(z) / U_{\infty}=0.14$ at $\mathrm{Q}(12.0, z)$. A flow reversal near the bed develops.
- In the plane $45^{\circ}$, the profiles notably in the upper region of the water depth decrease, having the maximum value below the water surface; but in the lower region there is tendency of increase.
- In the plane $90^{\circ}$, the profiles increase, attaining the maximum value of $\bar{u}(z) / U_{\infty}=1.7$ at $\mathrm{Q}(12.0, z)$. However, some points near the surface decrease, showing the negative values at points close to the cylinder. This is possibly a consequence of the variability of the separation point at the surface, observed visually during the measurements. This point moves upstream and downstream in plane $90^{\circ}$. The variation of the separation point is affected by the wake vortex behind the cylinder.
- In the plane $135^{\circ}$, the profiles first decrease, especially in the upper region of the water depth, forming a flow reversal, and then increase closer to the cylinder.
- In the plane $180^{\circ}$, the profiles decrease, forming a flow reversal in the upper region of the water depth.

Approaching the cylinder one observes for the transversal velocity profiles, $\bar{v}(z) / U_{\infty}$ the following:

- In the plane $0^{\circ}$, the profiles are rather uniform over the depth and their values are much smaller compared to the longitudinal ones.
- In the plane $45^{\circ}$, the profiles increase, having a rather constant value over the entire water depth, attaining the maximum value of $\bar{v}(z) / U_{\infty}=0.77$ at $\mathrm{Q}(12.0, z)$.
- In the plane $90^{\circ}$, the profiles first increase and then decrease, notably at the upper region of the water depth and near the bed. They have a concave distribution.
- In the plane $135^{\circ}$, the profiles increase at the upper region of the water depth, and decrease at the lower region.
- In the plane $180^{\circ}$, the profiles are chaotic.

Approaching the cylinder one observes for the vertical velocity profiles, $\bar{w}(z) / U_{\infty}$, as the following:

- In the plane $0^{\circ}$, the profiles increase, having always a downward direction. Its distributions have a zero value at the surface and a maximum value near the bed. The maximum downward flow is about $\bar{w}(z) / U_{\infty}=0.35$ at $\mathrm{Q}(12.0, z)$, at elevation near the bed.
- In the pl ane $45^{\circ}$, the profiles increase, having always a downward direction, and being rather constant values from the surface to near the bed.
- In the plane $90^{\circ}$, the profiles increase, having always a downward direction. The maximum downward velocity near the surface, as being $\bar{w}(z) / U_{\infty}=0.47$ at $\mathrm{Q}(12.0, z)$.
- In the plane $135^{\circ}$, the profiles increase, having downward direction.
- In the plane $180^{\circ}$, the profiles at $\mathrm{r}>27.5$ [cm] increase, having upward direction, and decrease at $r<27.5$ [cm], having downward direction.


Figure 1. Velocity distributions, $\bar{u}(z), \bar{v}(z), \bar{w}(z)$, measured in different planes around the cylinder, for Test 2

## Horizontal distributions of Velocities

The depth-averaged velocities, $U$ and $V$, obtained by the integration of the point velocity over the water depth, using the following definition:

$$
\begin{equation*}
U=1 / h \int \bar{u}(z) d z \quad \text { and } \quad V=1 / h \int \bar{v}(z) d z \tag{1}
\end{equation*}
$$

These are presented in the first pictures of Figure 2. The shadow plane behind the cylinder is the region of the wake, calculated by the numerical model at steady state condition. Results of the depth-averaged velocity, found from the numerical simulation ( see Yulistiyanto et.al,1997) are also plotted. The figures show the flow pattern around the cylinder. Approaching the cylinder, in plane $0^{\circ}$, the depthaveraged velocity decrease, however in plane $45^{\circ}$ and $90^{\circ}$, they increase, having the maximum values close to the cylinder. At the plane $180^{\circ}$, they show a velocity reversal at points near the cylinder and positive velocities at points far from the cylinder. The flow patterns in this plane, together with the plane $135^{\circ}$, show the existence of the wake behind the cylinder.

The point velocity at some relative elevations, at $z / h=0.96, z / h=0.5$, and $z / h=0.03$, are also presented in Figs. 2.

At $z / h=0.96$, one observes that in plane $0^{\circ}$ and $45^{\circ}$, the flow patterns are similar to the ones of the depth-averaged velocity. However in plane $180^{\circ}$, all point velocities are in the reversal direction. The flow patterns at this plane, together with the ones in plane $90^{\circ}$ and $135^{\circ}$, show the existence of the wake at the surface. The size of regions affected by the wake is bigger in Test 1 than in Test 2. This indicates that the larger flow Reynolds number has a larger wake.

At $z / h=0.5$, one observes that the flow pattern as well as its magnitudes are similar as the ones of the depth-averaged velocity.

At the relative elevation $z / h=0.03$, one observes that the flow patterns are different from the ones of the depth-averaged velocity. At the plane $0^{\circ}$, $45^{\circ}$, and $90^{\circ}$, this flow pattern is formed by the existence of the horseshoe vortex in plane $0^{\circ}$ which is then washed away downstream in plane $45^{\circ}$ and $90^{\circ}$.

## Velocity vectors in planes around the cylinder

Further interesting observations can be made, if the velocity vectors are shown in the different vertical planes, radiating around the cylinder. The above (see Figure 1) velocity profiles, $u(z)$ and $v(z)$, can be used to obtain the radial velocity, such as:

$$
\begin{equation*}
\bar{u}_{r}(z)=\bar{u} \cos (180-\alpha)+\bar{v} \sin (180-\alpha) \tag{2}
\end{equation*}
$$

shown also in the insert of Figure 4.
The vertical distributions of the velocities vectors, $\bar{u}$, and $\bar{w}$, have been plotted in Figure 3 for test 2; the same is shown but in a three-dimensional view in Figure 4 for test 1 . The following is to be observed:

- In the plane $\alpha=0^{\circ}$ : unidirectional flow is towards the cylinder diminishing in intensity and getting more and more two-directional, first in the upper layers and subsequently over the entire flow depth; a vertical downward flow in front of the cylinder is noticed and close to the base of the cylinder a flow reversal in the lowest layers is then evident.
- In the plane $\alpha=45^{\circ}$ : the same as in the plane $\alpha=$ $0^{\circ}$ is observed, but everything gets more pronounced, notably the flow reversal and the downward flow in front of the cylinder.
- In the plane $\alpha=90^{\circ}$ : flow is away from the cylinder, diminishing in intensity and getting more and more unidirectional; very close to the cylinder and over the entire depth a flow reversal is evident, notably in test 1.
- In the plane $\alpha=135^{\circ}$ : flow is in the zone of separation and the wake of the cylinder manifests itself; a strong flow reversal is seen closer to the cylinder and only far from the cylinder the flow becomes gradually unidirectional; in the lower layers the flow away is strong, while in the upper layers the flow is directed towards the cylinder.
- In the plane $\alpha=180^{\circ}$ : the flow pattern is rather similar to the one observed for the plane $\alpha=135^{\circ}$.

As can be seen in the planes of approach, $\left(\alpha=0^{\circ}\right.$ and $\alpha=45^{\circ}$, the flow moves into a region of an adverse pressure gradient, created by the cylinder. Consequently the bottom layers start to separate and climb up along a shear layer, beginning at the point of separation, $S_{v}$ (see Figure 3). Underneath this shear layer, in the bottom corner of the cylinder, a clockwise vortex is formed. Both incoming flow and the vortex are deflected sideways as a result of blockage by further incoming flow; they move downstream, with the axis of the vortex in the streamwise direction (see Figure 3 and Figure 4). In addition, the vertical downward flow close by the cylinder, augmenting the formation of clockwise vortex, begins to form a counter-flow, also referred as mushroom vortex (see Agui et Andreopoulos, 1992). What is seen here is a horseshoe-vortex system, developing in the bottom corner of the cylinder, confined to a flow depth of $z / h<0.15$.
In the planes behind the cylinder, $\alpha=135^{\circ}$ and $\alpha=$ $180^{\circ}$, the wake created by the cylinder is evident over
the entire flow depth. Flow far downstream of the
cylinder becomes again unidirectional.


Figure 2. Depth averaged and point velocities vectors of the flow around the cylinder


Figure 3. Velocity vectors in planes $0^{\circ}, 45^{\circ}, 90^{\circ}, 135^{\circ}$ and $180^{\circ}$ for Test 2


Figure 5. Scheme of the horseshoe-vortex system

## CONCLUSION

A schematic representation, given with Fig. 5, shall serve as a summary of our most important findings:

- Measurements of the instantaneous velocity upstream and around a cylinder are reported for fully developed turbulent open-channel flow at $R e_{D} \approx 10^{5}$. This three-dimensional flow has been documented with the mean point velocity field. The measurements were made with a nonintrusive ADVP-instrument.
- Due to the presence of the cylinder, the threedimensional separated flow is composed of a complex flow system, known as the horseshoevortex system. It is formed at the base of the cylinder and loops around the cylinder in the downstream direction.
- The vertical distributions of velocities of the flow around a cylinder show a highly threedimensional picture, being reasonably organized in the front and less organized in the back of the cylinder, where separations flow takes place.
- From the horizontal distributions of velocities, it can be seen the existence of the horizontal wake vortex behind the cylinder. The size of regions affected by this wake is bigger in flow with larger flow Reynolds number.
- Downstream of the cylinder, the wake created by the cylinder is evident over the entire flow depth.


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